SEARCHES FOR MONOPOLES, SUPERSYMMETRY, TECHNICOLOR, COMPOSITENESS, EXTRA DIMENSIONS, etc.

Magnetic Monopole Searches

Isolated supermassive monopole candidate events have not been confirmed. The most sensitive experiments obtain negative results.

Best cosmic-ray supermassive monopole flux limit:

$$<~1.4\times10^{-16}~{\rm cm^{-2}sr^{-1}s^{-1}}~~{\rm for}~1.1\times10^{-4}<\beta<1$$

Supersymmetric Particle Searches

All supersymmetric mass bounds here are model dependent.

The limits assume:

1) $\widetilde{\chi}_1^0$ is the lightest supersymmetric particle; 2) *R*-parity is conserved; See the Particle Listings for a Note giving details of supersymmetry.

$$\begin{array}{l} \widetilde{\chi}_i^0 \ -- \ \text{neutralinos} \ (\text{mixtures of} \ \widetilde{\gamma}, \ \widetilde{Z}^0, \ \text{and} \ \widetilde{H}_i^0) \\ \text{Mass} \ m_{\widetilde{\chi}_1^0} \ > \ 0 \ \text{GeV}, \ \text{CL} = 95\% \\ \text{[general MSSM, non-universal gaugino masses]} \\ \text{Mass} \ m_{\widetilde{\chi}_1^0} \ > \ 46 \ \text{GeV}, \ \text{CL} = 95\% \\ \text{[all } \tan\beta, \ \text{all} \ m_0, \ \text{all} \ m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}] \\ \text{Mass} \ m_{\widetilde{\chi}_2^0} \ > \ 620 \ \text{GeV}, \ \text{CL} = 95\% \\ \text{[}\widetilde{\chi}_2^0 \ \to \ \ell^\pm \ell^\mp \widetilde{\chi}_1^0, \ \text{simplified model}, \ m_{\widetilde{\chi}_1^0} = 0 \ \text{GeV}] \\ \text{Mass} \ m_{\widetilde{\chi}_3^0} \ > \ 620 \ \text{GeV}, \ \text{CL} = 95\% \\ \text{[}\widetilde{\chi}_3^0 \ \to \ \ell^\pm \ell^\mp \widetilde{\chi}_1^0, \ \text{simplified model}, \ m_{\widetilde{\chi}_1^0} = 0 \ \text{GeV}] \\ \text{Mass} \ m_{\widetilde{\chi}_4^0} \ > \ 116 \ \text{GeV}, \ \text{CL} = 95\% \\ \text{[}1 < \tan\beta < 40, \ \text{all} \ m_0, \ \text{all} \ m_{\widetilde{\chi}_2^0} \ - m_{\widetilde{\chi}_1^0}] \\ \end{array}$$

$$\begin{split} \widetilde{\chi}_i^{\pm} &- \text{charginos (mixtures of } \widetilde{W}^{\pm} \text{ and } \widetilde{H}_i^{\pm}) \\ &\text{Mass } m_{\widetilde{\chi}_1^{\pm}} > 94 \text{ GeV, CL} = 95\% \\ &\text{[} \tan\beta < 40, \ m_{\widetilde{\chi}_1^{\pm}} - m_{\widetilde{\chi}_1^0} > 3 \text{ GeV, all } m_0] \\ &\text{Mass } m_{\widetilde{\chi}_1^{\pm}} > 500 \text{ GeV, CL} = 95\% \\ &\text{[simplified model, } 2\ell^{\pm} + E_T, \ m_{\widetilde{\chi}_1^0} = 0 \text{ GeV}] \\ \widetilde{\chi}^{\pm} &- \text{long-lived chargino} \\ &\text{Mass } m_{\widetilde{\chi}^{\pm}} > 620 \text{ GeV, CL} = 95\% \quad \text{[stable } \widetilde{\chi}^{\pm}] \\ \widetilde{\nu} &- \text{sneutrino} \\ &\text{Mass } m > 41 \text{ GeV, CL} = 95\% \quad \text{[model independent]} \\ &\text{Mass } m > 94 \text{ GeV, CL} = 95\% \quad \text{[CMSSM, } 1 \leq \tan\beta \leq 40, \ m_{\widetilde{e}_R} - m_{\widetilde{\chi}_1^0} > 10 \text{ GeV}] \\ \widetilde{e} &- \text{scalar electron (selectron)} \\ &\text{Mass } m(\widetilde{e}_L) > 107 \text{ GeV, CL} = 95\% \quad \text{[all } m_{\widetilde{e}_L} - m_{\widetilde{\chi}_1^0}] \\ \widetilde{\mu} &- \text{scalar muon (smuon)} \\ &\text{Mass } m > 94 \text{ GeV, CL} = 95\% \quad \text{[CMSSM, } 1 \leq \tan\beta \leq 40, \ m_{\widetilde{\mu}_R} - m_{\widetilde{\chi}_1^0} > 10 \text{ GeV}] \\ \widetilde{\tau} &- \text{scalar tau (stau)} \end{split}$$

Where the limits below show a **range** of lower bounds, the bounds depend on different simplified models, different signals, different assumptions, and different luminosities.

 $\begin{array}{ll} [m_{\widetilde{\tau}_R}-m_{\widetilde{\chi}_1^0}>& 15 \text{ GeV, all } \theta_\tau, \text{ B}(\widetilde{\tau}\to~\tau\,\widetilde{\chi}_1^0)=100\%] \\ \text{Mass } m>~286 \text{ GeV, CL}=95\% \quad [\text{long-lived }\widetilde{\tau}] \end{array}$

 \widetilde{q} – squarks of the first two quark generations

Mass m > 81.9 GeV, CL = 95%

The first of these limits is within CMSSM with cascade decays, evaluated assuming a fixed value of the parameters μ and $\tan\beta$, and assuming two-generations of mass degenerate squarks $(\widetilde{q}_L \text{ and } \widetilde{q}_R)$ and gaugino mass parameters that are constrained by the unification condition at the grand unification scale.

Mass
$$m>1450$$
 GeV, CL = 95%
 [CMSSM, $\tan\beta=30$, $A_0=-2\max(m_0,\,m_{1/2})$, $\mu>0$] Mass $m>608-1260$ GeV, CL = 95%
 [mass degenerate squarks] Mass $m>490-600$ GeV, CL = 95%
 [single light squark bounds]

 \widetilde{q} — long-lived squark

Mass m>1000, CL = 95% \widetilde{t} , charge-suppressed interaction model]

Mass m>845, CL = 95% \widetilde{b} , stable, Regge model] \widetilde{b} — scalar bottom (sbottom)

Mass m>323-880 GeV, CL = 95% \widetilde{t} — scalar top (stop)

Mass m>323-800 GeV, CL = 95%

[Lower value is a decay via charm, and upper a decay via top] \widetilde{g} — gluino

Mass m>700-1780 GeV, CL = 95%

Technicolor

The limits for technicolor (and top-color) particles are quite varied depending on assumptions. See the Technicolor section of the full *Review* (the data listings).

Quark and Lepton Compositeness, Searches for

Scale Limits Λ for Contact Interactions (the lowest dimensional interactions with four fermions)

If the Lagrangian has the form

$$\pm \frac{g^2}{2\Lambda^2} \overline{\psi}_{\mathsf{L}} \gamma_{\mu} \psi_{\mathsf{L}} \overline{\psi}_{\mathsf{L}} \gamma^{\mu} \psi_{\mathsf{L}}$$

(with $g^2/4\pi$ set equal to 1), then we define $\Lambda \equiv \Lambda_{LL}^{\pm}$. For the full definitions and for other forms, see the Note in the Listings on Searches for Quark and Lepton Compositeness in the full *Review* and the original literature.

$$\Lambda_{LL}^{+}(eeee)$$
 > 8.3 TeV, CL = 95%
 $\Lambda_{LL}^{-}(eeee)$ > 10.3 TeV, CL = 95%
 $\Lambda_{LL}^{+}(ee\mu\mu)$ > 8.5 TeV, CL = 95%
 $\Lambda_{LL}^{-}(ee\mu\mu)$ > 9.5 TeV, CL = 95%
 $\Lambda_{LL}^{+}(ee\tau\tau)$ > 7.9 TeV, CL = 95%

$$\begin{array}{lll} \Lambda_{LL}^{-}(e\,e\,\tau\,\tau) &> 7.2 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{+}(\ell\ell\ell\ell) &> 9.1 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(\ell\ell\ell\ell) &> 10.3 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{+}(e\,e\,q\,q) &> 16.4 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,q\,q) &> 20.7 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,u\,u) &> 23.3 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,u\,u) &> 12.5 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,d\,d) &> 11.1 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,d\,d) &> 26.4 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,c\,c) &> 9.4 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,b\,b) &> 9.4 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(e\,e\,b\,b) &> 10.2 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(\mu\mu q\,q) &> 15.8 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(\mu\mu q\,q) &> 21.8 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(\mu\mu q\,q) &> 2.81 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(q\,q\,q\,q) &> 12.0 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(q\,q\,q\,q) &> 17.5 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(\mu\nu q\,q) &> 5.0 \text{ TeV, CL} = 95\% \\ \Lambda_{LL}^{-}(\nu\nu q\,q) &> 5.4 \text{ TeV, CL}$$

Excited Leptons

The limits from $\ell^{*+}\ell^{*-}$ do not depend on λ (where λ is the $\ell\ell^{*}$ transition coupling). The λ -dependent limits assume chiral coupling.

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\begin{array}{l} e^{*\pm} - \text{excited electron} \\ \text{Mass } m > \ 103.2 \text{ GeV}, \ \text{CL} = 95\% \quad (\text{from } e^* \, e^*) \\ \text{Mass } m > \ 3.000 \times 10^3 \text{ GeV}, \ \text{CL} = 95\% \quad (\text{from } e \, e^*) \\ \text{Mass } m > \ 356 \text{ GeV}, \ \text{CL} = 95\% \quad (\text{if } \lambda_{\gamma} = 1) \\ \\ \mu^{*\pm} - \text{excited muon} \\ \text{Mass } m > \ 103.2 \text{ GeV}, \ \text{CL} = 95\% \quad (\text{from } \mu^* \mu^*) \\ \text{Mass } m > \ 3.000 \times 10^3 \text{ GeV}, \ \text{CL} = 95\% \quad (\text{from } \mu \mu^*) \\ \end{array}
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 $\tau^{*\pm} - \text{excited tau} \\ \text{Mass } m > 103.2 \text{ GeV, CL} = 95\% \quad (\text{from } \tau^*\tau^*) \\ \text{Mass } m > 2.500 \times 10^3 \text{ GeV, CL} = 95\% \quad (\text{from } \tau\tau^*) \\ \nu^* - \text{excited neutrino} \\ \text{Mass } m > 1.600 \times 10^3 \text{ GeV, CL} = 95\% \quad (\text{from } \nu^*\nu^*) \\ \text{Mass } m > 213 \text{ GeV, CL} = 95\% \quad (\text{from } \nu^*X) \\ q^* - \text{excited quark} \\ \text{Mass } m > 338 \text{ GeV, CL} = 95\% \quad (\text{from } q^*q^*) \\ \text{Mass } m > 5.200 \times 10^3 \text{ GeV, CL} = 95\% \quad (\text{from } q^*X) \\ \end{aligned}$

Color Sextet and Octet Particles

Color Sextet Quarks (q_6) Mass m>84 GeV, CL=95% (Stable q_6)
Color Octet Charged Leptons (ℓ_8) Mass m>86 GeV, CL=95% (Stable ℓ_8)
Color Octet Neutrinos (ν_8) Mass m>110 GeV, CL=90% $(\nu_8\to \nu_g)$

Extra Dimensions

Please refer to the Extra Dimensions section of the full *Review* for a discussion of the model-dependence of these bounds, and further constraints.

Constraints on the radius of the extra dimensions, for the case of two-flat dimensions of equal radii

 $R < 30~\mu\text{m}$, CL = 95% (direct tests of Newton's law) $R < 10.9~\mu\text{m}$, CL = 95% ($p\,p \to j\,G$) R < 0.16–916 nm (astrophysics; limits depend on technique and assumptions)

Constraints on the fundamental gravity scale

$$M_{TT}>$$
 6.3 TeV, CL $=95\%$ $~~$ (pp $\rightarrow~$ dijet, angular distribution) $M_c>$ 4.16 TeV, CL $=95\%$ $~~$ (pp $\rightarrow~$ $\ell \overline{\ell})$

Constraints on the Kaluza-Klein graviton in warped extra dimensions

$$M_G > 3.3$$
 TeV, CL = 95% $(pp \rightarrow e^+e^-, \mu^+\mu^-)$

Constraints on the Kaluza-Klein gluon in warped extra dimensions

$$M_{g_{KK}}~>~2.5$$
 TeV, CL $=95\%~~(g_{KK}
ightarrow~t\,\overline{t})$